



difference equations, one finds a regularization method for computing such representations by solving a linear system of equations [5]. By solving such a linear system, the issue of discretizing kernels near its singularity is avoided.

In many problems, the choice of basis should accommodate not only the integral operators but also differential operators and the boundary conditions. Multiwavelet bases developed in [6] satisfy many of these requirements. As is well known, multiwavelet bases retain some properties of wavelet bases, such as vanishing moments, orthogonality, and compact support. The basis functions do not overlap on a given scale and are organized in small groups of several functions (thus, multiwavelets) sharing the same support. On the other hand, some of the basis functions are discontinuous, similar to the Haar basis and in contrast to wavelets with regularity. Because of the vanishing moments of the basis functions, a wide class of integro-differential operators has effectively sparse representations in these bases. (By an effectively sparse matrix representation, we mean a representation that differs from a sparse matrix by a matrix with a small norm.) More recently, it was demonstrated in [7] that such bases are useful for adaptively solving partial differential equations (PDEs) with boundary conditions.

As a part of the program to develop multidimensional adaptive PDE solvers, we construct representations for homogeneous convolution operators in dimensions $d = 2, 3$, and higher. Another method using representations of low separation rank for functions and operators is described in [2], and we outline this approach as well. Let us start with the straightforward multiwavelet generalization of [1]. For the multiwavelet bases, the computational costs are $O(k^4 N)$ in two dimensions and $O(k^6 N)$ in three dimensions, where order- k multiwavelets are used and N is the number of boxes in which significant coefficients exist. In many applications, such as computational chemistry, it is too expensive to compute with such algorithms. The algorithms that we present here have computational complexity of $O(k^2 N)$ and $O(k^3 N)$, respectively.

In our solution method, the use of a localized, discontinuous, and adaptive basis of multiwavelets is combined with a representation of functions and operators that generalizes the separation of variables. This is all performed with controlled accuracy in finite-precision arithmetic. The notable features of multiwavelet representation of operators are the following:

- Multiwavelets form an orthonormal basis.

Although it is not known whether all operators and functions in practical applications have a short LSR representation, many important operators, such as the multiparticle Schrödinger operator and the inverse Laplacian, can be efficiently represented in this form.

We apply our framework for numerical calculus of operators. These operations are important in applications in which functions of operators must be computed. For example, the Schultz iteration¹ for the computation of the inverse requires operator products and sums.

Let us start by providing a formal description of a multiresolution analysis (MRA) for a multiwavelet [11, 12]. Such an MRA is defined as an ascending chain of embedded closed subspaces of the Hilbert space $L_2([0, 1])$,

$$\dots \subset V_0 \subset V_1 \subset V_2 \subset \dots,$$

with the properties that

$$L_2([0, 1]) = \overline{\cup_j V_j}, \quad \cap_j V_j = \{0\}.$$

Additional requirements are as follows:

1. The subspace V_0 is invariant under integer translations.
2. The subspaces V_j are all scaled version of one another. One or more scaling functions are in V_j if and only if $(2^j x)$ is in V_0 .
3. One or more scaling functions are in V_0 such that their rescaled and shifted versions, of the form $2^{j/2} \phi(2^j x - k)$, constitute an orthonormal basis² of V_j .

Well-known examples are the Haar basis, in which the scaling function is the characteristic function restricted to the interval $[0, 1]$, the Battle-Lemarie wavelet with spline scaling function, and the Daubechies families of wavelets [4], as well as multiwavelets [6]. F6sw3inri6.0801 39scend], asc 7encew3in

in the multiwavelet basis, the function f is expressed as

$$f(x) = \sum_j s_j \phi_j^0(x) + \sum_{n=0}^{n-1} \sum_{l=0}^{2^n-1} \sum_{j=0}^{k-1} d_{jl}^{(n)} \phi_{jl}^{(n)}(x).$$

The scalars $s_{jl}^{(n)} = \int f(x) \phi_{jl}^{(n)}(x) dx$ and $d_{jl}^{(n)} = \int f(x) \phi_{jl}^{(n)}(x) dx$ are respectively called the scaling and multiwavelet coefficients. These can be computed directly using Gauss-Legendre quadrature or by the following two-scale relations from a fine scale to a coarse scale.

The scaling and multiwavelet functions satisfy the two-scale relations

$$\phi_j(x) = \sqrt{2} \sum_{j=0}^{k-1} [h_{ij}^{(0)} \phi_j(2x) + h_{ij}^{(1)} \phi_j(2x-1)]$$

and

$$\phi_j(x) = \sqrt{2} \sum_{j=0}^{k-1} [g_{ij}^{(0)} \phi_j(2x) + g_{ij}^{(1)} \phi_j(2x-1)],$$

where $h_{ij}^{(0)}$, $h_{ij}^{(1)}$, $g_{ij}^{(0)}$, and $g_{ij}^{(1)}$ are coefficients which are easily computed given the scaling and multiwavelet basis. By using these relations, further two-scale relations can be derived for the scaling and multiwavelet coefficients from scale n and $n+1$,

$$s_{ij}^{(n)} = \sqrt{2} \sum_{j=0}^{k-1} [h_{ij}^{(0)} s_{j,2l}^{(n+1)} + h_{ij}^{(1)} s_{j,2l+1}^{(n+1)}]$$

and

$$d_{ij}^{(n)} = \sqrt{2} \sum_{j=0}^{k-1} [g_{ij}^{(0)} s_{j,2l}^{(n+1)} + g_{ij}^{(1)} s_{j,2l+1}^{(n+1)}]$$

of the basis functions. For example, in two dimensions, the
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1. We assume that the scaling function coefficients $[r_l]_{ij}$ are known for a sufficiently large distance $|l|$ from the singularity.
2. The two-scale relations (e.g., for $[r_l]_{ij}$) are used to compute $[r_l]_{ij}$ for $1 < |l| < n$.
3. The two-scale relations produce a system of linear equations for coefficients in the range $|l| \leq 1$. Coefficients in this range appear on both sides of the two-scale difference equations.

In general, this criterion is written as

$$2^{n+} r = Ar + b,$$

where r represents the vector of coefficients. The matrix A consists of combinations of the coefficients of two-scale relations.

Let $x, y \in R^n$ and let T be a convolution operator with the kernel $K(x, y)$, homogeneous of degree $-n$. Assuming that the solution to the linear system in condition 3 exists, we obtain a multiresolution kernel $T_0(x, y)$ with coefficients r_l from the construction above. We define the multiresolution regularized operators to be $T_j : V_j \rightarrow V_j, j \in Z$, with kernels $T_j(x, y) = 2^{-j-n} T_0(2^{-j}x, 2^{-j}y)$ on the chosen MRA as $j \rightarrow \infty$.

We illustrate this construction in one dimension. The two-scale relations for the coefficients representing the kernel with homogeneity degree $-n$ are

$$[r_l]_{ij} = 2^{-n-} \sum_{l',j'=0}^{k-1} h_{il'}^{(0)} h_{jj'}^{(1)} [r_{2l-1}]_{l'j'} + h_{il}^{(0)} h_{jj}^{(0)} - h_{il}^{(1)} h_{jj}^{(1)} [r_{2l}]_{l'j'} - h_{il}^{(1)} h_{jj}^{(1)} [r_{2l+1}]_{l'j'},$$

If the kernel has singularity at the origin and the multiwavelet coefficients are known for $|l| > 1$, the two-scale relations consist of three equations:

~~$$[r_0]_{ij} = 2^{-n-} \sum_{l',j'=0}^{k-1} h_{il'}^{(0)} h_{jj'}^{(1)} [r_{-1}]_{l'j'} + h_{il}^{(0)} h_{jj}^{(0)} - h_{il}^{(1)} h_{jj}^{(1)} [r_0]_{l'j'} - h_{il}^{(1)} h_{jj}^{(1)} [r_1]_{l'j'}$$

$$2^{-n-}$$~~

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In [3] a multiwavelet method is applied to solve the Schrödinger equation:

$$\left(-\frac{1}{2}\Delta + V\right)\Psi = E\Psi.$$

The integral form of this equation in three dimensions is

$$\Psi(r) = \frac{1}{4}$$

new $O(N)$ multiscale simulation capabilities. The sparse representations of these operators were used to produce multiresolution methods for applying the Hilbert transform and solving the Poisson and Schrödinger

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G B Fa *Oak Ridge National Laboratory, Oak Ridge, Tennessee 37831 (gif@ornl.gov)*. Dr. Fann is a Staff Member at the Oak Ridge National Laboratory. He received an A.B. degree from the University of California at Berkeley, and M.S. and Ph.D. degrees from MIT, all in mathematics. He received an R&D 100 Prize in 1999 for the parallel eigensolver PeIGS used in the computational chemistry software NWChem and other computational chemistry software while a staff member of the Pacific Northwest National Laboratory. He was the architect of the parallel imaging software PeICEIS. Dr. Fann's interests include fast computational methods and parallel computing. He is the author of more than 30 publications.

G B B *Department of Applied Mathematics, University of Colorado, Boulder, Colorado 80309 (beylkin@colorado.edu)*. Dr. Beylkin is a Professor of Applied Mathematics at the University of Colorado at Boulder. He received a Diploma in mathematics from the University of St. Petersburg (Leningrad), FSU, in 1975 and a Ph.D. in mathematics from the Courant Institute of Mathematical Sciences, NYU, in 1982. After one year as a Postdoctoral Fellow at NYU, he joined Schlumberger-Doll Research in Ridgefield, Connecticut, where he worked on wave propagation, signal processing, and inverse problems. He spent the year 1988-1989 at Yale University working on fast algorithms and applications of wavelets. In 1991 he joined the faculty of the Program in Applied Mathematics (later the Department of Applied Mathematics) at the University of Colorado at Boulder. Professor Beylkin's interests include fast numerical algorithms, approximation theory, wave propagation, and inverse problems. He has published more than 60 papers and holds one U.S. patent. He is a member of AMS, SIAM, and SEG.

RSb J. Ha B *University of Tennessee, Knoxville, and Oak Ridge National Laboratory, Oak Ridge, Tennessee 37831 (harrisonrj@ornl.gov)*. Dr. Harrison received a Ph.D. degree from Cambridge University, England, in 1984. He currently works at the Oak Ridge National Laboratory, with a joint research appointment at the University of Tennessee. His major research interests are the electronic structure theory of molecules with an emphasis on high precision, and high-performance computing. Dr. Harrison is the main architect of NWChem, a large computational chemistry package for parallel computers for which he received an R&D 100 award in 1999 and the IEEE Sydney Fernbach Award in 2002. He is the author of more than 70 publications.

K E. JB da *IBM Life Sciences Solutions, One Rogers Street, Cambridge, Massachusetts 02142 (kjordan@us.ibm.com)*. Dr. Jordan is Emerging Solution Executive in IBM Life Sciences Solutions and team leader for the Life Sciences Strategic Relationships and Institutes of Innovation Programs. In this role, he has responsibility for overseeing Life Sciences university partnerships, investigating and developing concepts for new areas of growth for IBM in the life sciences, especially involving high-performance computing, and providing leadership in high-end computing and simulation in areas such as systems biology and high-end visualization. He received a B.S. degree in mathematics from Hobart College, and M.S. and Ph.D. degrees in applied mathematics from the University of Delaware, the latter in 1980. Prior to joining IBM, he held positions at Exxon Research and Engineering,

Thinking Machines, and Kendall Square Research. He is currently Vice President for Industry in the Society for Industrial and Applied Mathematics (SIAM). Dr. Jordan has received several awards for his work on supercomputers, and his main research interests are in the efficient use of advanced architecture computers for modeling and simulation. He is an author of more than 30 papers on subjects that include interactive visualization on parallel computers, parallel domain decomposition for reservoir/groundwater simulation, turbulent convection flows, parallel spectral methods, multigrid techniques, multiresolution wavelets, and wave propagation.